

# Un invito al Trasporto Ottimo Quantistico<sup>1</sup>

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<sup>1</sup>Based on joint works with G. De Palma, M. Marvian, S. Lloyd, T. Titkos and D. Viosztek

# Outline

- 1 Classical Optimal Transport
- 2 Quantum Systems
- 3 Quantum Optimal Transport
- 4 Conclusion

# Plan

- 1 Classical Optimal Transport
  - Monge
  - Kantorovich
  - Earth Mover's distance
- 2 Quantum Systems
- 3 Quantum Optimal Transport
- 4 Conclusion

# Monge's transport problem

Monge (1781): *sur la théorie des déblais et des remblais*.



How to **transport** soil during a construction with **minimal expenses**?

# The assignment problem

A discrete formulation: given a

- cost  $c(x, y)$  of moving unit of soil from position  $x$  to position  $y$ , e.g.

$$c(x, y) = |x - y|,$$

- Source distribution of soil  $\sigma = (\sigma(x_i))_i$
- Target distribution (dump)  $\rho = (\rho(y_j))_j$

Find  $T : \{x_i\} \rightarrow \{y_j\}$  that moves  $\sigma$  into  $\rho$  with **minimal transport cost**

$$\sum_i c(x_i, T(x_i))\sigma(x_i).$$

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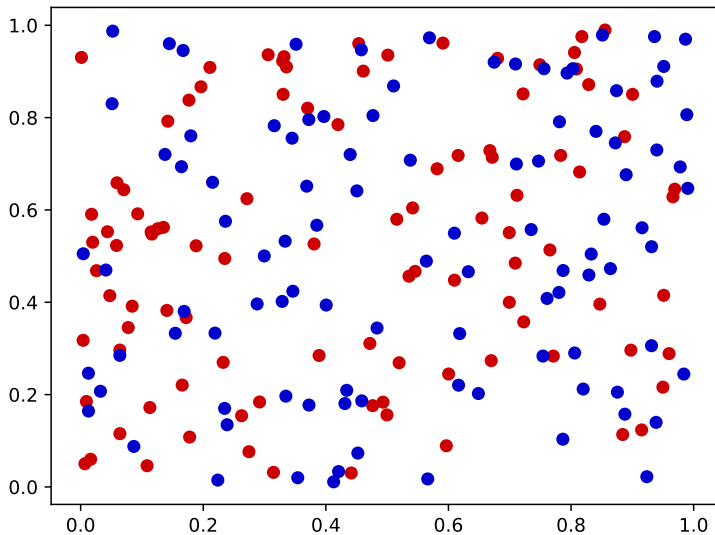
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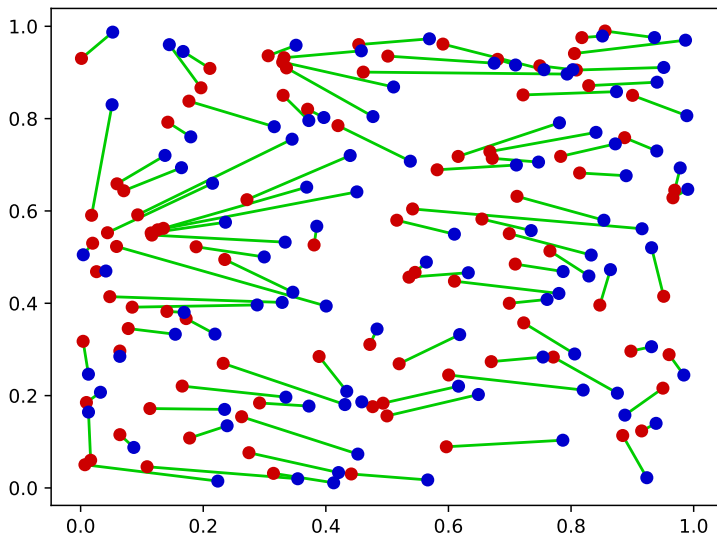
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# Kantorovich and linear programming

Relax the map  $T$  to a **transport plan**

$$T(x_i, y_j) \geq 0$$

such that

$$\sum_j T(x_i, y_j) = \sigma(x_i), \quad \sum_i T(x_i, y_j) = \rho(y_j).$$

Probabilistic interpretation:

$$K(y_j|x_i) := \frac{T(x_i, y_j)}{\sigma(x_i)} \in [0, 1].$$

The variational problem becomes

$$\min_T \sum_i \sum_j c(x_i, y_j) T(x_i, y_j)$$

⇒ **linear programming!**

# From soil to probabilities

- For a transport plan  $T$  it must be

$$\sum_i \sigma(x_i) = \sum_i \sum_j T(x_i, y_j) = \sum_j \sum_i T(x_i, y_j) = \sum_j \rho(y_j).$$

- We assume that  $\rho, \sigma$  are probability mass functions (discrete densities):

$$\sum_i \sigma(x_i) = \sum_j \rho(y_j) = 1.$$

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# Earth Mover's distance and duality

- If  $c(x, y) = d(x, y)$  is a distance, then

$$W_1(\sigma, \rho) = \min_T \sum_i \sum_j d(x_i, y_j) T(x_i, y_j)$$

induces a distance (Wasserstein-Kantorovich, Earth Mover's).

- Kantorovich duality:

$$W_1(\sigma, \rho) = \max \left\{ \sum_i f(x_i) \sigma(x_i) - \sum_j f(y_j) \rho(y_j) : |f(x) - f(y)| \leq d(x, y) \right\}.$$

- Buy at price  $f(y_j)$  and sell at price  $f(x_i) \Rightarrow$  maximize the profit!

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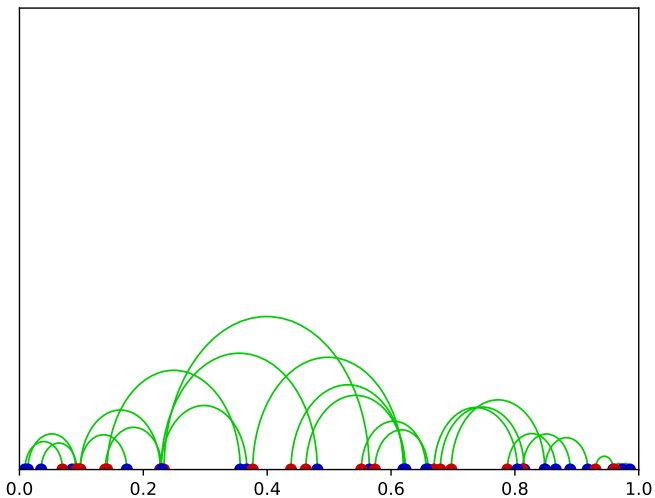
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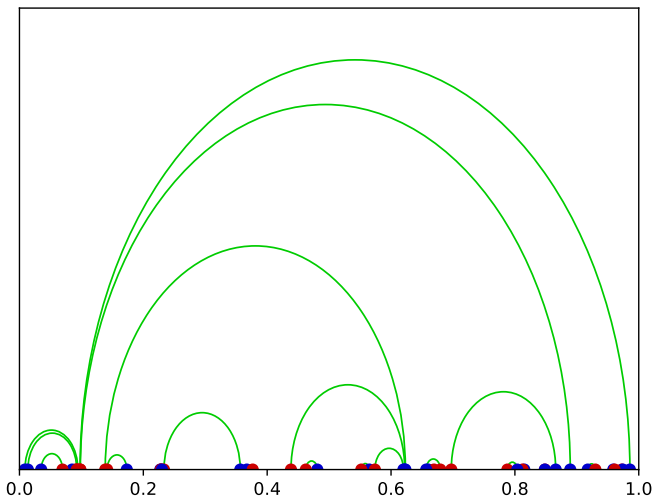
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$$c(x, y) = |x - y|^p, p = 1$$

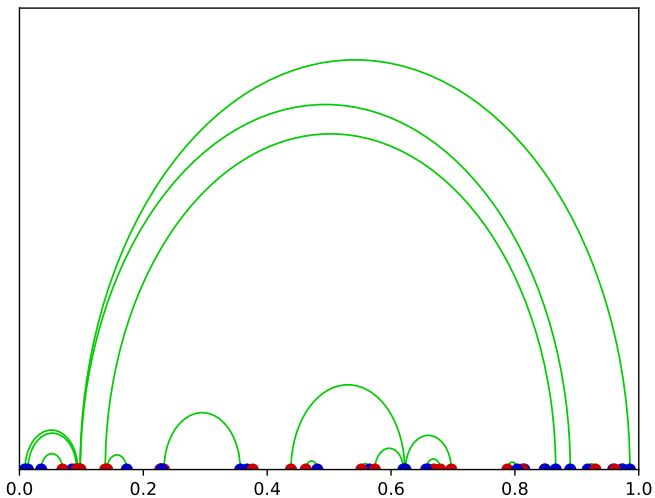




$$c(x, y) = |x - y|^p, p = 0.6$$



$$c(x, y) = |x - y|^p, p = 0.1$$



Some applications of optimal transport:

- 1 comparison between **point clouds**
- 2 geometric **interpolation** between distributions
- 3 discriminator in **generative AI models** (WGANs)
- 4 functional inequalities (isoperimetric, concentration of measure)
- 5 PDE's as **gradient flows**
- 6 geometry (synthetic **Ricci curvature bounds**)

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# Plan

- 1 Classical Optimal Transport
- 2 Quantum Systems**
  - From Classical to Quantum
  - Systems of qubits
- 3 Quantum Optimal Transport
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# Classical vs Quantum: a dictionary

Classical



Quantum



# Classical vs Quantum: a dictionary

---

$E$  (finite set)

$e \in E$

$A \subseteq E$

---

$f : E \rightarrow \mathbb{C}$

real-valued

non-negative

$|f|^2$

---

$\sum_{x \in E} f(x)$

---

$\mathcal{H}$  Hilbert space ( $\mathbb{C}^n$ )

$|\psi\rangle \in \mathcal{H}$

$\mathcal{S} \subseteq \mathcal{H}$

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$A : \mathcal{H} \rightarrow \mathcal{H}$  linear

self-adjoint ( $A = A^\dagger$ )

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$A^2 = A A$

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probabilities  $p \in \mathcal{P}(E)$

Dirac  $p = \delta_x$

Markov (transition) operator

$K : \mathcal{P}(E) \rightarrow \mathcal{P}(E)$

$(Kp)(y) = \sum_{x \in E} K(y|x)p(x)$

Shannon's entropy

$S(p) = - \sum_x p(x) \log p(x)$

Relative entropy (KL divergence)

$D(p||q) = \sum_x p(x) \log(p(x)/q(x))$

Wavefunction  $\psi \in \mathcal{S}(H)$

Pure state  $\rho = |\psi\rangle\langle\psi|$

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# Single qubit system

A quantum analogue of  $\{0, 1\}$ . Set

$$H = \mathbb{C}^2.$$

Standard (**computational**) basis

$$\{|0\rangle, |1\rangle\} = \{(1, 0), (0, 1)\},$$

Pauli operators

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Parametrization of states:

$$\rho = \frac{1}{2} (\mathbb{I}_{\mathbb{C}^2} + b_x \sigma_x + b_y \sigma_y + b_z \sigma_z),$$

with

$$(b_x, b_y, b_z) \in \mathbb{R}^3, \quad b_x^2 + b_y^2 + b_z^2 \leq 1.$$

$$\rho = |\psi\rangle\langle\psi| \text{ is pure} \Leftrightarrow b_x^2 + b_y^2 + b_z^2 = 1.$$

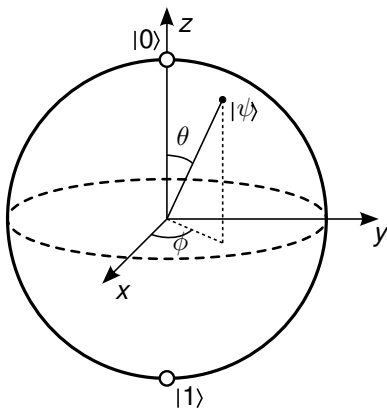


Figure: The Bloch representation of a pure state  $\rho = |\psi\rangle\langle\psi| \in \mathbb{C}^2$ .

## Many-qubits systems

A quantum analogue of  $\{0, 1\}^n$ :

$$H_n = (\mathbb{C}^2)^{\otimes n} \sim \mathbb{C}^{2^n}.$$

Standard (computational) basis

$$\{|s\rangle : s \in \{0, 1\}^n\}.$$

**Pure states** are  $\rho = |\psi\rangle\langle\psi|$  but not necessarily  $\psi \in \{0, 1\}^n$ :

$$|\psi\rangle = \frac{1}{2^{n/2}} \sum_{s \in \{0,1\}^n} |s\rangle \quad (\text{uniform superposition})$$

$$|\Phi^+\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \quad (\text{Bell state, } n = 2)$$

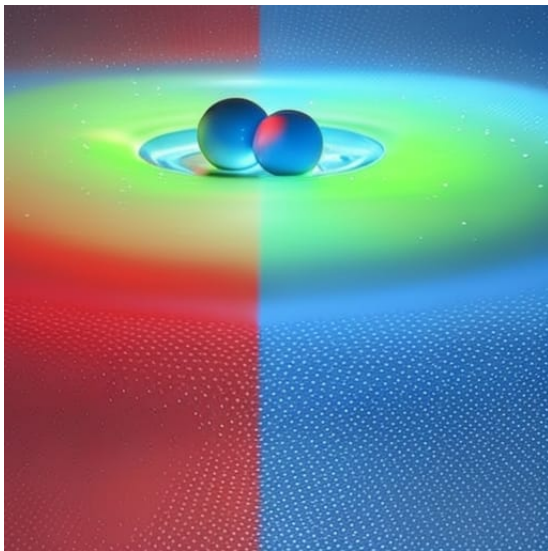
**“Classical” probabilities** on  $\{0, 1\}^n$  are diagonal in the computational basis:

$$\rho = \sum_{s \in \{0,1\}^n} p(s) |s\rangle\langle s|.$$

# Plan

- 1 Classical Optimal Transport
- 2 Quantum Systems
- 3 Quantum Optimal Transport**
  - Overview
  - Transport via quantum channels
  - Quantum earth mover's distance
- 4 Conclusion

# Why Quantum Optimal Transport?



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Classical distances between probabilities have quantum analogues:

- Total variation  $\rightarrow$  Trace distance
- Hellinger distance  $\rightarrow$  Fidelity
- KL divergence  $\rightarrow$  Relative entropy

Like their classical counterparts:

- They generally apply to comparisons of distributions
- They are not to be confused with quantum distance

What about Quantum Optimal Transport?

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• They are not so easy to compute quantumly

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Like their classical counterparts:

- Quite general, easy to compute or approximate
- Easy to extend to quantum probability

What about Quantum Optimal Transport?



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What about **Quantum Optimal Transport**?

# A timeline

2007 - Villani, OT

optimal transport in non-convexity geometry

2007 - Ambrosio, Gigli, Savaré

Wasserstein distance of curved manifolds

2008 - Ambrosio, Gigli

quantum analogue of Brenier's Kantor formula

2009 - Ambrosio

Wasserstein extending any distance on Lipschitz

2010 - Ambrosio, Gigli

quantum Kantorovich's problem

2010 - Ambrosio, Gigli

Wasserstein distance on metric spaces

2010 - Ambrosio, Gigli, Savaré

Wasserstein distance on Lie groups

# A timeline

- 1992 – Connes/Lott:  
spectral distance in non-commutative geometry
- 1997 – Ziemer/Blumenthal:  
Wasserstein distance of Poisson distributions
- 2002 – Bapat/Prati:  
quantum analogue of Jensen-Shannon formula
- 2012 – Ambrose:  
1-Wasserstein extending any distance on finite vectors
- 2015 – O’Connell/Carlen/Truitt:  
quantum Kantorovich’s problem
- 2016 – De Palma/Truitt:  
quantum Wasserstein distance
- 2017 – De Palma/Truitt/Li:  
quantum Wasserstein distance

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- 1997 - Zyczkowski/Slomczynski:  
Wasserstein distance of Husimi distributions
- 2012 - Maas/Carlen:  
quantum analogue of Benamou-Brenier formula
- 2013 - Agredo:  
1-Wasserstein extending any distance on basis vectors
- 2016 - Golse/Mouhot/Paul:  
quantum Kantorovich problem
- 2019 - De Palma/T.:  
Quantum optimal transport using channels
- 2020 - De Palma/Marvian/T./Lloyd:  
Earth mover's distance on qubits



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# Transport via quantum channels

- Any Kantorovich transport plan  $T(x_i, y_j)$  yields a transition kernel

$$K(y_j|x_i) = \frac{T(x_i, y_j)}{\sigma(x_i)},$$

such that

$$K\sigma = \rho$$

- Fix real-valued functions  $(g_\ell)_{\ell=1}^d$  and choose a “quadratic” cost

$$c(x, y) = \sum_{\ell=1}^d (g_\ell(x) - g_\ell(y))^2.$$

$$\begin{aligned} & \sum_i \sum_j c(x_i, y_j) T(x_i, y_j) \\ &= \sum_{\ell} \sum_i g_\ell^2(x_i) \sigma(x_i) + \sum_j g_\ell^2(y_j) \rho(y_j) - 2 \sum_i g_\ell(x_i) (K^\dagger g_\ell)(x_i) \sigma(x_i), \end{aligned}$$

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$$(K^\dagger g)(x_i) = \sum_j g(y_j) K(y_j|x_i)$$

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- Given  $\rho, \sigma \in \mathcal{S}(H)$ , **define** transport plans as quantum channels

$$\Phi : \mathcal{S}(H) \rightarrow \mathcal{S}(H), \quad \Phi(\tau) = \sum_j B_j \tau B_j^\dagger$$

such that

$$\Phi(\sigma) = \rho.$$

- Fix observables  $\mathcal{R} = (R_\ell)_{\ell=1}^d$  and **define**

$$C_{\mathcal{R}}(\Phi) = \sum_{\ell} \text{Tr}[R_\ell^2 \sigma] + \text{Tr}[R_\ell^2 \rho] - 2 \text{Tr}[R_\ell \sqrt{\sigma} (\Phi^\dagger R_\ell) \sqrt{\sigma}],$$

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**Problem:** Describe the isometries  $J : \mathcal{S}(H) \rightarrow \mathcal{S}(H)$

$$D_{\mathcal{R}}(J(\rho), J(\sigma)) = D_{\mathcal{R}}(\rho, \sigma) \quad \forall \rho, \sigma.$$

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# Quantum earth mover's distance

- **Aim:** define a quantum analogue of earth mover's distance with respect to the **Hamming** distance on binary strings:

$$|x - y|_H = \sum_{i=1}^n |x_i - y_i| \quad \text{for } x, y \in \{0, 1\}^n.$$

- Recall the Kantorovich duality:

$$W_1(\sigma, \rho) = \max \left\{ \sum_{x \in \{0,1\}^n} f(x)(\sigma(x) - \rho(x)) : |f(x) - f(y)| \leq |x - y|_H \right\}.$$

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# Quantum Lipschitz constant

- Classical  $L$ -Lipschitz functions with respect to the Hamming distance have **bounded differences**:

$$|f(x) - f(y)| \leq L \quad \forall x, y \in \{0, 1\}^n \text{ differing in only one site.}$$

- Equivalently:

$$\|f\|_{Lip} = 2 \max_{i=1, \dots, n} \min \{ \|f - g_i\|_{\infty} : g_i : \{0, 1\}^n \rightarrow \mathbb{R} \text{ independent of site } i \}$$

- Given an observable  $A$  on  $H_n = (\mathbb{C}^2)^{\otimes n}$ , define

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$$\|f\|_{Lip} = 2 \max_{i=1, \dots, n} \min \{ \|f - g_i\|_{\infty} : g_i : \{0, 1\}^n \rightarrow \mathbb{R} \text{ independent of site } i \}$$

- Given an observable  $A$  on  $H_n = (\mathbb{C}^2)^{\otimes n}$ , **define**

$$\|A\|_{Lip} = 2 \max_{i=1, \dots, n} \min \{ \|A - A_i \otimes \mathbb{I}_{\mathbb{C}^2}\|_{\infty} : A_i \text{ observable on } H_{n-1} \}.$$

- We **define** the quantum earth mover's distance by duality:

$$W_1(\sigma, \rho) = \max \{ \text{Tr}[A(\sigma - \rho)] : \|A\|_{Lip} \leq 1 \}.$$

# Gaussian concentration

- Let  $\sigma = 2^{-n} \mathbb{I}_{(\mathbb{C}^2)^{\otimes n}}$  (corresponding to a classical uniform distribution)
- Quantum Marton's inequality: for any state  $\rho$ ,

$$W_1(\sigma, \rho) \leq \sqrt{\frac{n}{2}} S(\rho \| \sigma).$$

- Duality  $\Rightarrow$  if  $A$  satisfies  $\text{Tr}[A] = 0$  and  $\|A\|_{Lip} \leq 1$ , then

$$\dim(A \geq k\sqrt{n}\mathbb{I}_{(\mathbb{C}^2)^{\otimes n}}) \leq 2^n e^{-2k^2}.$$

- Informally:

$$A \approx 2^{-n} \text{Tr}[A] \pm \|A\|_{Lip} \sqrt{n}.$$

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# Plan

- 1 Classical Optimal Transport
- 2 Quantum Systems
- 3 Quantum Optimal Transport
- 4 Conclusion**

# Perspectives





# Perspectives

Quantum OT has already found applications:

- Non-commutative **Ricci curvature** bounds
- Quantum machine learning (generative models, quantum neural nets)
- Equilibrium measures on quantum spin systems

... many others just around the corner?

Further mathematical developments:

- $\mathbb{R}$ -valued Ricci curvature (Bakry-Ellis, Otto)
- Geometry induced by  $\mathcal{L}_1$  (Bernot, Indricato)
- Entropy and  $\mathcal{L}_1$  (Bakry, Pratelli)
- Quantum transportation (Baccelli, Caracciolo)

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• **Quantum Ricci curvature** (Ollivier, Villani)

• **Geometry induced by OT** (Gromov, Ambrosio, Pratelli)

• **Non-commutative OT** (Gheorghiu)

• **Quantum Wasserstein** (Gheorghiu, Ollivier, Villani)

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• **Quantum Ricci curvature**

• **Quantum entropy measures** (e.g. **quantum Rényi entropy**)

• **Quantum Wasserstein distances**

• **Quantum Pinsker inequality**

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